

# Lecture 9, Feb 5, 2024

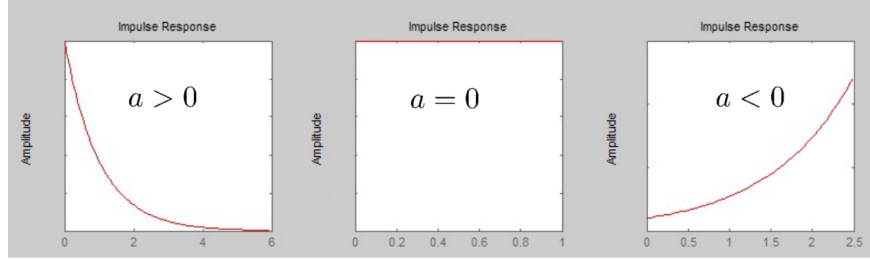
## First-Order System Response

- Consider a pure integrator:  $y(t) = \int_0^t u(t) dt + y(0)$  which has transfer function  $H(s) = \frac{1}{s}$  if  $y(0) = 0$ 
  - The ODE is  $\dot{y}(t) = u(t)$
  - The impulse response is  $y_i(t) = \mathcal{L}^{-1}\{H(s)\} = \mathcal{L}^{-1}\left\{\frac{1}{s}\right\} = 1$
  - The step response is  $y_s(t) = \mathcal{L}^{-1}\left\{H(s)\frac{1}{s}\right\} = \mathcal{L}^{-1}\left\{\frac{1}{s^2}\right\} = t$
  - What if the initial condition is not zero?
    - \* Laplace transform the ODE to get  $sY(s) - y(0) = U(s) \implies Y(s) = \frac{1}{s}U(s) + \frac{1}{s}y(0)$
    - \* For a step response,  $y_s(t) = \mathcal{L}^{-1}\left\{\frac{1}{s^2}\right\} + \mathcal{L}^{-1}\left\{\frac{1}{s}y(0)\right\} = t + y(0)1(t) = t + y(0)$
- Consider an RC circuit with input voltage  $Ku(t)$ 
  - Form the ODE:  $T\dot{y}(t) + y(t) = Ku(t)$  where  $T = RC$
  - Laplace transform:  $TsY(s) + Y(s) = KU(s) \implies H(s) = \frac{Y(s)}{U(s)} = \frac{K}{Ts + 1}$
  - Impulse response:  $y_i(t) = \mathcal{L}^{-1}\{H(s)\} = \frac{K}{T}e^{-\frac{t}{T}}$
  - Step response:  $y_s(t) = \mathcal{L}^{-1}\left\{\frac{K}{s(Ts + 1)}\right\} = K\mathcal{L}^{-1}\left\{\frac{1}{s} - \frac{T}{Ts + 1}\right\} = K\left(1 - e^{-\frac{t}{T}}\right)$
  - We can see that  $T$  is the time constant of the system; the smaller it is, the faster the system evolves
  - DC gain:  $y_s s = \lim_{t \rightarrow \infty} y_s(t) = \lim_{s \rightarrow 0} s \frac{K}{s(Ts + 1)} = K$
- In general a first-order system has transfer function  $H(s) = \frac{b}{s + a}$  and impulse response  $h(t) = be^{-at}1(t)$ 
  - For positive  $a$ , this is stable and the system decays to 0; for negative  $a$ , this is unstable; for  $a = 0$  the system maintains a constant output
    - \* Positive  $a$  gives poles in the LHP and negative  $a$  gives poles in the RHP
  - The step response is given by  $y_s = \frac{b}{a}(1 - e^{-at})1(t)$ 
    - \* For positive  $a$ , this converges to the DC gain  $\frac{b}{a}$
    - \* For negative  $a$  this diverges exponentially
    - \* For zero  $a$  this gives a linear response (note we can derive this by noting  $H(s) = \frac{b}{s}$  in this case)
- The time constant is given by  $T = \frac{1}{a}$ 
  - The rise time is given by  $t_r \approx 2.2T$ , which is the time taken for the output to go from 10% to 90% of the DC gain
  - The settling time is given by  $t_s \approx \frac{4.6}{a}$ , which is the time taken for the output to reach 99% of the DC gain
- In a first-order system, there is never any overshoot or oscillation; the output never passes the steady state value

## Second Order System Response

- Consider a spring-mass-dashpot system:  $m\ddot{y}(t) + b\dot{y}(t) + ky(t) = kf(t)$ 
  - Laplace transform:  $m(s^2Y(s) - sy(0^-) - \dot{y}(0^-)) + b(sY(s) - y(0^-)) + kY(s) = kF(s)$
  - $Y(s) = \frac{\frac{k}{m}}{s^2 + \frac{b}{m}s + \frac{k}{m}}F(s) + \frac{s + \frac{b}{m}}{s^2 + \frac{b}{m}s + \frac{k}{m}}y(0^-) + \frac{1}{s^2 + \frac{b}{m}s + \frac{k}{m}}\dot{y}(0^-)$

➤ Impulse Response:  $h(t) = \mathcal{L}^{-1}[H(s)] = b e^{-at} 1(t)$



➤ Step Response:  $y_{step}(t) = \mathcal{L}^{-1}[H(s)/s] = \frac{b}{a}(1 - e^{-at})1(t)$

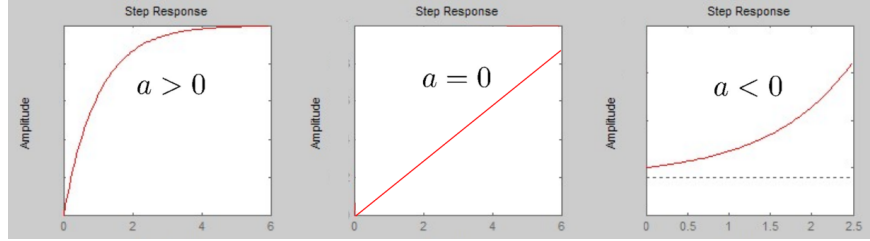


Figure 1: Behaviour of the impulse and step responses for a general (strictly proper) first-order system.

- Assuming zero state,  $H(s) = \frac{Y(s)}{F(s)} = \frac{\frac{k}{m}}{s^2 + \frac{b}{m}s + \frac{k}{m}} = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2}$ 
  - \*  $\omega_n = \sqrt{\frac{k}{m}}$  is the natural frequency
  - \*  $\zeta = \frac{b}{2\sqrt{km}}$  is the damping ratio
- The poles are at  $-\zeta\omega_n \pm \omega_n\sqrt{\zeta^2 - 1}$ 
  - \* Depending on  $\zeta$  we can get real or imaginary poles
- If  $\zeta > 1$  (i.e.  $b > 2\sqrt{km}$ ) we have two distinct real poles; the system is *overdamped*
  - \* Let  $-\sigma_1 = -\zeta\omega_n + \omega_n\sqrt{\zeta^2 - 1}$ ,  $-\sigma_2 = -\zeta\omega_n - \omega_n\sqrt{\zeta^2 - 1}$
  - \* Then  $\omega_n = \sqrt{\sigma_1\sigma_2}$ ,  $\zeta = \frac{\sigma_1 + \sigma_2}{2\sqrt{\sigma_1\sigma_2}}$
- If  $\zeta = 1$  (i.e.  $b = 2\sqrt{km}$ ) we have two overlapping real poles; the system is *critically damped*
  - \*  $H(s) = \frac{\sigma^2}{(s + \sigma)^2}$  where  $\sigma = \omega_n$
- If  $0 \leq \zeta < 1$  (i.e.  $b < 2\sqrt{km}$ ) we have two complex conjugate poles; the system is *underdamped*
  - \* The poles are  $s_1, s_2 = -\sigma \pm j\omega_d$  where  $\sigma = \zeta\omega_n$ ,  $\omega_d = \omega_n\sqrt{1 - \zeta^2}$
  - \*  $H(s) = \frac{\omega_n^2}{(s - (-\sigma + j\omega_d))(s - (-\sigma - j\omega_d))} = \frac{\sigma^2 + \omega_d^2}{(s + \sigma)^2 + \omega_d^2}$
  - \* In this case the system oscillates
  - \*  $\omega_d$  is the oscillation frequency and  $\sigma$  is the decay rate
- Consider the impulse response of the underdamped case
  - $y_i(t) = \mathcal{L}^{-1} \left\{ \frac{(\sigma^2 + \omega_d^2)}{(s + \sigma)^2 + \omega_d^2} \right\} = \mathcal{L}^{-1} \left\{ \frac{(\sigma^2 + \omega_d^2)}{\omega_d} \frac{\omega_d}{(s + \sigma)^2 + \omega_d^2} \right\} = \frac{\sigma^2 + \omega_d^2}{\omega_d} e^{-\sigma t} \sin(\omega_d t)$
  - Alternatively  $y_i(t) = \frac{\omega_n}{\sqrt{1 - \zeta^2}} e^{-\zeta\omega_n t} \sin(\omega_n \sqrt{1 - \zeta^2} t)$
  - The response is a decaying exponential

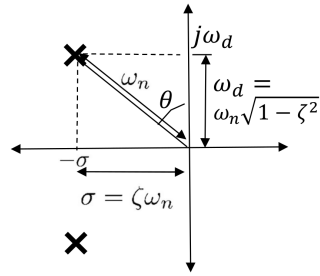


Figure 2: Illustration of the system variables in polar form.

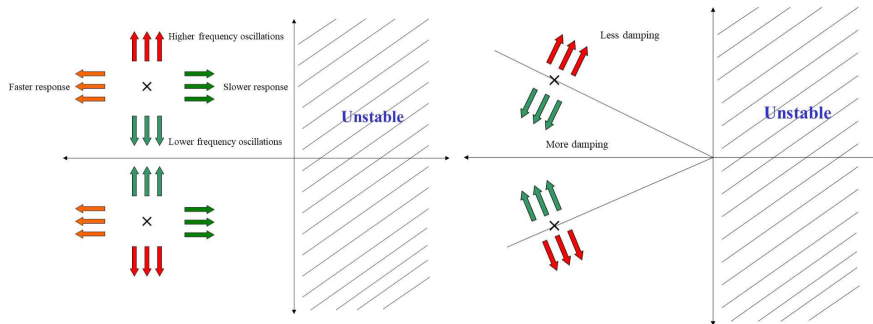


Figure 3: Response of an underdamped second-order system based on pole location.

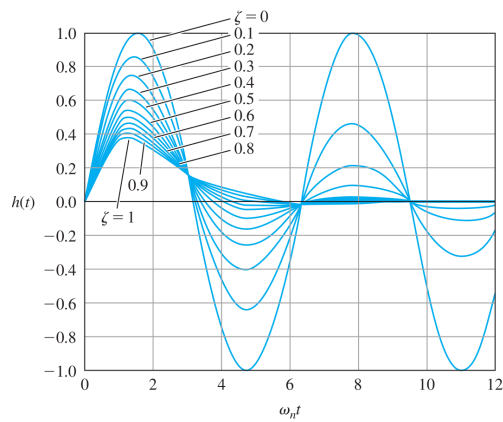


Figure 4: Impulse response of an underdamped second-order system.