

## Lecture 28 (2-11), Mar 23, 2023

### Ideal Gas (Sackur-Tetrode Formula)

- For a given  $U$ , how many microstates are there?
- $U = \sum_{i=1}^N \frac{p_i^2}{2m} \implies \sum_{i=1}^N \vec{p}_i^2 = \sum_{i=1}^N p_{xi}^2 + p_{yi}^2 + p_{zi}^2 = 2mU$ 
  - How many ways can be distribute this momentum to get the same  $U$ ?
  - We want the sum of the squares of  $3N$  numbers to equal  $2mU$
  - This is a sphere that lives in  $3N$ -dimensional space, or  $S^{3N-1}$  in  $\mathbb{R}^{3N}$ , with radius  $r^2 = 2mU \implies r = \sqrt{2mU}$
- $p_x, p_y, p_z$  are all quantized in units of  $\frac{\pi\hbar}{L}$  (where  $L^3 = V$ ), like in the case of the infinite square well; these points form a grid in space, the sphere will hit some of these grid points, and each hit is a microstate
  - We only need the sphere to be close to these grid points
  - We expect that the number of hits is proportional to the area of the sphere, proportional to  $r^{3N-1}$ , with  $r$  defined by the energy
- The area of a  $3N - 1$ -dimensional sphere in  $3N$  dimensions can be shown to be  $\frac{2\pi^{\frac{3N}{2}}}{\Gamma(\frac{3N}{2})} R^{3N-1}$ 
  - If  $\frac{3N}{2}$  is an integer this is  $\frac{2\pi^{\frac{3N}{2}}}{(\frac{3N}{2} - 1)!}$
  - Therefore  $\Omega(N, U) \sim \frac{2\pi^{\frac{3}{2}N}}{(\frac{3}{2}N - 1)!} (\sqrt{2mU})^{3N-1}$ , but we have to fix it first:
    - \* The dimensions aren't consistent: to fix it, we divide  $\sqrt{2mU}$  by a factor of  $\frac{\pi\hbar}{L}$  first to make it unitless and make the units match
      - This basically converts the  $p_i$  to  $n_i$
    - \* The  $n_i$  are always positive, so the  $p_i$  must all be positive; this means we have to reduce the surface area by a factor of 2 for every axis
      - e.g. for a circle, if we restrict it to  $x, y > 0$ , we have to divide by a factor of 4; for a sphere, restricting it to the first quadrant divides by a factor of 8
    - \* The particles should be indistinguishable according to quantum mechanics, which means that if we swap the momenta of two particles, it stays in the same microstate
      - There are  $N!$  ways to permute the momenta, which all lead to identical microstates, so we have to reduce  $\Omega$  by a factor of  $N!$
      - This is called the Gibbs factor, which he derived before QM
  - $\Omega(N, U) = \frac{2\pi^{\frac{3}{2}N}}{N! (\frac{3}{2}N - 1)! 2^{3N}} \left( \frac{\sqrt{2mUL}}{\pi\hbar} \right)^{3N-1}$
- Now to find the entropy  $S = k \ln \Omega$ 
  - First use Stirling's approximation,  $N! = \left(\frac{N}{e}\right)^N$  and ignore the -1
  - Substitute  $L = V^{\frac{1}{3}}$
  - $\Omega(N, U) = \frac{e^N}{N^N} \frac{2}{(2^3)^N} \frac{2(\pi^{\frac{3}{2}})^N (e^{\frac{3}{2}})^N}{(3^{\frac{3}{2}})^N (N^{\frac{3}{2}})^N} \frac{V^N}{(\pi\hbar)^{3N}} \left( (2mU)^{\frac{3}{2}} \right)^N$
  - = ... I gave up ...
  - $S = k \ln \Omega$ 
    - =  $k \ln 2 + kN \ln \left( \left( \left( \frac{4\pi mU}{N} \right)^{\frac{3}{2}} \frac{1}{(2\pi\hbar)^3} \frac{V}{N} \right) + \frac{5}{2} \right)$
    - =  $kN \ln \left( \left( \left( \frac{4\pi mU}{N} \right)^{\frac{3}{2}} \frac{1}{(2\pi\hbar)^3} \frac{V}{N} \right) + \frac{5}{2} \right)$

\* This is known as the Sackur-Tetrode formula for the entropy of an ideal gas

- Taking  $\frac{1}{T} = \frac{\partial S}{\partial U} = Nk \frac{3}{2} \frac{1}{U} \implies U = \frac{3}{2} kNT$  which is the same as the one we got from equipartition
- What can we learn from this formula?
  - $S$  is on the same order as  $kN$
  - For all the extensive quantities, e.g.  $N, U, V$ , if we double them, the entropy also doubles
  - $S$  increases with  $U$ , and levels off – this is a normal system, the temperature increases with energy and we have a positive heat capacity
  - If we lower  $U$  enough, then eventually the stuff inside the log will become less than 1, and we'll get a negative entropy, which is impossible – therefore this formula does not work for gases with very low energy