

## Lecture 26, Mar 20, 2023

### Magnetic Dipole

- A *magnetic dipole* is simply a closed loop of current, characterized by its *magnetic dipole moment*  $\vec{m} = IS\hat{a}_n$ , where the direction is determined through the right hand rule and  $S$  is the enclosed area
  - e.g. for a loop with radius  $a$ , we have  $\vec{m} = \pi a^2 I \hat{a}_z$
  - If the loop has  $n$  turns, then the effective  $I$  is increased, so the magnetic dipole moment is magnified by a factor of  $n$
  - $\vec{m} = nIS\hat{a}_n$  with units  $[\text{A m}^2]$
  - A magnetic dipole will produce a field in the same direction as the direction it points in
- What happens to a magnetic dipole moment in a  $\vec{B}$  field?
  - The loop will experience some magnetic force  $\vec{F} = I\vec{L} \times \vec{B}$  (from  $\vec{F}_m = q\vec{u} \times \vec{B}$ )
  - This produces a net torque  $\vec{T} = \vec{m} \times \vec{B}$
  - When  $\vec{m}$  and  $\vec{B}$  are aligned, the torque goes to zero; therefore a magnetic dipole will rotate until its own field is aligned with the applied field

### Magnetization

- All materials have small atomic magnetic dipoles caused by the movement of electrons around the nuclei
  - Since they're all randomly oriented, there is no net field
- In a magnetic material, in the presence of an external magnetic field, the dipoles experience a torque that aligns them in the same direction as the field
- The overall result is that the small  $\vec{B}$  fields from the dipoles now all point in the same direction, producing a net magnetic field
  - The magnetic field produced by the dipoles is in the same direction as the external applied field, so they add together
- A material is *magnetic* if it allows their atomic magnetic dipoles to be all aligned in the same direction
- Define the *magnetization vector*  $\vec{M}$  (akin to  $\vec{P}$ 's relationship with  $\vec{p}$ ) as an average of the magnetic dipoles within a material:
  - $\vec{M} = \lim_{\Delta v \rightarrow 0} \frac{1}{\Delta v} \sum_i \vec{m}_i \approx N\vec{m}$  with units  $[\text{A/m}]$
- This magnetization leads to a *bound current density* (surface)  $\vec{J}_{ms} = \vec{M} \times \hat{a}_n$  where  $\hat{a}_n$  is the outward normal of the surface
  - There could also be volume bound current densities
- Now we can define 3 new quantities:
  - The magnetic field intensity  $\vec{H} = \frac{\vec{B}}{\mu_0} - \vec{M} = \frac{\vec{B}}{\mu_r \mu_0}$
  - The magnetic susceptibility  $\chi_m$ , where  $\vec{M} = \chi_m \vec{H}$
  - The relative permeability  $\mu_r = \chi_m + 1$
  - Like in the electric field case,  $\vec{B}$  accounts for both bound and free currents, but  $\vec{H}$  only cares about free currents

#### Important

The magnetic flux density in a magnetized material is not always greater than the applied field, since the magnetic dipole moments can also align to be antiparallel to the applied field, depending on the material (as a consequence,  $\chi_m$  is not necessarily positive, so  $\mu_r$  could be less than 1)

- Example: cylindrical permanent magnet, where a constant uniform  $M = M_0 \hat{a}_z$  exists; the cylinder is defined by  $-\frac{L}{2} \leq z \leq \frac{L}{2}, 0 \leq r \leq a$ 
  - $\vec{J}_{ms} = \vec{M} \times \hat{a}_n = \vec{M} \times \hat{a}_r = M_0 \hat{a}_\phi$

$$\begin{aligned}
- \vec{B} &= \iint \frac{\mu_0 \vec{J}_{ms} \times (\vec{R} - \vec{R}')}{4\pi |\vec{R} - \vec{R}'|^3} a d\phi' dz' \\
&= \frac{\mu_0}{4\pi} \int_{-\frac{L}{2}}^{\frac{L}{2}} \int_0^{2\pi} \frac{M_0 \hat{a}_\phi \times (-a\hat{a}_r + (z - z')\hat{a}_z)}{(a^2 + (z - z')^2)^{\frac{3}{2}}} a d\phi' dz' \\
&= \frac{\mu_0 M_0}{2} \left( \frac{\frac{L}{2}}{\sqrt{a^2 + (z - L/2)^2}} + \frac{\frac{L}{2}}{\sqrt{a^2 + (z + L/2)^2}} \right) \hat{a}_z \\
- \text{What if } L \gg a? \\
&* \vec{B} \rightarrow \mu_0 M_0 \hat{a}_z
\end{aligned}$$

### Summary

When a magnetic material is exposed to an external applied magnetic field, it is magnetized; the magnetization is characterized by the magnetization vector,

$$\vec{M} = \chi_m \vec{H}$$

where the magnetic field intensity is defined as

$$\vec{H} = \frac{\vec{B}}{\mu_0} - \vec{M} = \frac{\vec{B}}{\mu_r \mu_0}$$

where the relative permeability is defined as

$$\mu_r = \chi_m + 1$$

The magnetization creates a surface bound current density,

$$\vec{J}_{ms} = \vec{M} \times \hat{a}_n$$

where  $\hat{a}_n$  is the surface normal vector, and also a volume bound current density,

$$\vec{J}_m = \vec{\nabla} \times \vec{M}$$