Lecture 5, Sep 21, 2023

Transport Equations

- We want to know how v is related to v° $-v = \mathcal{F}_a^T v_a = \mathcal{F}_b^T v_b$ $\underbrace{ \vec{v}}_{a} = \underbrace{\vec{v}}_{a} \underbrace{ \vec{v}}_{a} \underbrace{ \vec{v}}_{b} \underbrace{ \vec{v}}_{b} \underbrace{ \vec{F}}_{a} \underbrace{ \vec{V}}_{b} \underbrace{ \vec{V}}_{b} \underbrace{ \vec{F}}_{a} \underbrace{ \vec{V}}_{b} \underbrace$ the scalar reference frames does not matter
- Therefore <u>v</u> = <u>v</u> [°] + <u>ω</u>^{ba} × <u>v</u>, which is known as the *transport equation* for velocity
 For acceleration, <u>v</u> = (<u>v</u> [°] + <u>ω</u>^{ba} × <u>v</u>).

$$= (\underline{v}^{\circ})^{\cdot} \underline{\omega}^{ba^{\cdot}} \times \underline{v} + \underline{\omega}^{ba} \times \underline{v}^{\cdot}$$

$$= \underline{v}^{\circ\circ} + \underline{\omega}^{ba} \times \underline{v}^{\circ} (\underline{\omega}^{ba^{\circ}} + \underline{\omega}^{ba} \times \underline{\omega}^{ba}) \times \underline{v} + \underline{\omega}^{ba} \times (\underline{v}^{\circ} + \underline{\omega}^{ba} \times \underline{v})$$

$$= \underline{v}^{\circ\circ} + 2\underline{\omega}^{ba} \times \underline{v}^{\circ} + \underline{\omega}^{ba^{\circ}} \times \underline{v} + \underline{\omega}^{ba} \times (\underline{\omega}^{ba} \times \underline{v})$$

- Notice that $\underline{\omega}^{ba} = \underline{\omega}^{ba}$
- If we interpret \vec{v} as a position vector, then $2\vec{\omega}^{ba} \times \vec{v}^{\circ}$ would be the Coriolis acceleration, $\vec{\omega}^{ba} \times (\vec{\omega}^{ba} \times \vec{r})$ would be the centripetal acceleration and $\vec{\omega}^{ba^{\circ}} \times \vec{r}$ would be the tangential acceleration (or Euler acceleration)
- In coordinate form, the velocity transport equation is $\dot{v}_a = C_{ab}(\dot{v}_b + {\omega_b^{ba}}^{\times} v_b)$
- For acceleration this is $\ddot{v}_a = C_{ab}(\ddot{v}_b + 2\omega_b^{ba^{\times}}\dot{v}_b + \dot{\omega}_b^{ba^{\times}}v + \omega_b^{ba^{\times}}\omega_b^{ba^{\times}}v_b)$ How do rotation matrices change with rotating reference frames?
- - $\mathcal{F}_{b}^{T} = \mathcal{F}_{a}^{T} \dot{C}_{ab} \text{ since } \mathcal{F}_{a}^{T} = \mathbf{0}$ $\omega_{b}^{ba^{\times}} = \mathcal{F}_{b} \cdot \mathcal{F}_{b}^{T} = C_{ba} \mathcal{F}_{a} \cdot \mathcal{F}_{a}^{T} \dot{C}_{ab} = C_{ba} \dot{C}_{ab}$ $\text{ If we transpose both sides we get } -\omega_{b}^{ba^{\times}} = \dot{C}_{ba} C_{ab}$

 - Multiplying by C_{ba} and rearranging, we get $\dot{C}_{ba} + \omega_b^{ba} \times C_{ba} = 0$
 - * We have found a differential equation for the rotation matrix that describes its evolution * This is known as Poisson's kinematical equation
- Consider now 3 reference frames $\mathcal{F}_a, \mathcal{F}_b, \mathcal{F}_c$; what is the relationship among $\underline{\omega}^{ba}, \underline{\omega}^{cb}$ and $\underline{\omega}^{ca}$?

$$-C_{ca} = C_{cb}C_{ba} + C_{cb}C_{ba}$$

$$-\omega_c^{ca\times} = -\dot{C}_{ca}C_{ac}$$

$$= -\dot{C}_{cb}C_{ba}C_{ab}C_{bc} - C_{cb}\dot{C}_{ba}C_{ab}C_{bc}$$

$$= -\dot{C}_{cb}C_{bc} - C_{cb}\dot{C}_{ba}C_{ab}C_{bc}$$

$$= \omega_c^{cb\times} + C_{cb}\omega_b^{ba\times}C_{bc}$$

$$= \omega_c^{cb\times} + (C_{cb}\omega_b^{ba})^{\times}$$

$$= (\omega_c^{cb} + C_{cb}\omega_b^{ba})^{\times}$$

- Therefore $\boldsymbol{\omega}_c^{ca} = \boldsymbol{\omega}_c^{cb} + \boldsymbol{C}_{cb} \boldsymbol{\omega}_b^{ba}$ - If we multiply both sides by $\boldsymbol{\mathcal{F}}_c^T$, we see that $\boldsymbol{\omega}^{ca} = \boldsymbol{\omega}^{cb} + \boldsymbol{\omega}^{ba}$

Important

While angular velocities can be added directly as $\underline{\omega}^{ca} = \underline{\omega}^{cb} + \underline{\omega}^{ba}$, angular accelerations cannot!

Summary

The transport equations relate velocity and acceleration as measured in one frame to how they're measured in another frame:

•
$$\vec{v} = \vec{y}^{\circ} + \vec{\omega}^{oa} \times \vec{v}$$

• $\vec{v}_a = C_{ab}(\vec{v}_b + \vec{\omega}_b^{ba} \times v_b)$

•
$$\vec{y} = \vec{y}^{\circ\circ} + 2\vec{\omega}^{oa} \times \vec{y}^{\circ} + \vec{\omega}^{oa} \times \vec{y} + \vec{\omega}^{oa} \times (\vec{\omega}^{oa} \times \vec{y})$$

• $\ddot{v}_a = \vec{C}_{ab}(\ddot{v}_b + 2\omega_b^{ba^{\times}}\dot{v}_b + \dot{\omega}_b^{ba^{\times}}\dot{v} + \omega_b^{ba^{\times}}\omega_b^{ba^{\times}}v_b)$

Rotation Representations Revisited

- Given an axis-angle representation a, ϕ , the Euler parameters are $\eta = \cos \frac{1}{2} \phi$ and $\varepsilon = a \sin \frac{1}{2} \phi$
 - η and ε are not independent, because we stipulate that $\eta^2 + \varepsilon^T \varepsilon = 1$
 - Euler parameters don't have a singularity
 - These are also known as quaternions
- Consider a 1-2-3 set of Euler angles, so $C = C_3 C_2 C_1$

$$\begin{aligned} - \omega^{\times} &= -CC^{T} \text{ (Note here } \omega = \omega_{b}^{0a}) \\ - \dot{C} &= C_{3}C_{2}\dot{C}_{1} + C_{3}\dot{C}_{2}C_{1} + \dot{C}_{3}C_{2}C_{1} \\ - \omega^{\times} &= -C_{3}C_{2}\dot{C}_{1}C_{1}^{T}C_{2}^{T}C_{3}^{T} - C_{3}\dot{C}_{2}C_{1}C_{1}^{T}C_{2}^{T}C_{3}^{T} - \dot{C}_{3}C_{2}C_{1}C_{1}^{T}C_{2}^{T}C_{3}^{T} \\ &= -C_{3}C_{2}\dot{C}_{1}C_{1}^{T}C_{2}^{T}C_{3}^{T} - C_{3}\dot{C}_{2}C_{2}^{T}C_{3}^{T} - \dot{C}_{3}C_{3}^{T} \\ &= C_{3}C_{2}\mathbf{1}_{1}^{\times}\dot{\theta}_{1}C_{2}^{T}C_{3}^{T} - C_{3}\mathbf{1}_{2}^{\times}\dot{\theta}_{2}C_{3}^{T} - \mathbf{1}_{3}^{\times}\dot{\theta}_{3} \\ &= (C_{3}C_{2}\mathbf{1}_{1}\dot{\theta}_{1})^{\times} + (C_{3}\mathbf{1}_{2}\dot{\theta}_{2})^{\times} + \mathbf{1}_{3}^{\times}\dot{\theta}_{3} \\ ^{*} \text{ Note that } \dot{C}_{1}C_{1}^{T} \text{ is the angular velocity of a rotation about the first axis, so it is equal to} \\ & \begin{bmatrix} 0 \\ 0 \\ \dot{\theta}_{1} \end{bmatrix}^{\times} = \mathbf{1}_{1}^{\times}\dot{\theta}_{1} \end{aligned}$$

- This means $\boldsymbol{\omega} = \boldsymbol{C}_3 \boldsymbol{C}_2 \boldsymbol{1}_1 \dot{\theta}_1 + \boldsymbol{C}_3 \boldsymbol{1}_2 \dot{\theta}_2 + \boldsymbol{1}_3 \dot{\theta}_3 = \begin{bmatrix} \boldsymbol{C}_3 \boldsymbol{C}_2 \boldsymbol{1}_1 & \boldsymbol{C}_3 \boldsymbol{1}_2 & \boldsymbol{1}_3 \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \\ \dot{\theta}_3 \end{bmatrix} = \boldsymbol{S}(\theta_2, \theta_3) \dot{\boldsymbol{\theta}}$

- * This is the kinematic equation for a 1-2-3 set of Euler angles (page 89 lists the S matrices for other combinations)
- At a singularity, S becomes singular (hence the name singularity) given ω , at a singularity, we cannot find $\dot{\theta}$
- Given any rotation representation, we can write a kinematic equation for it