## Lecture 8, Sep 26, 2022

## Spring-Coupled Masses

• Two equal spring oscillators coupled by another spring

• 
$$\begin{cases} m \frac{\mathrm{d}^2 x_A}{\mathrm{d}t^2} = -kx_A + k(x_B - x_A) = kx_B - 2kx_A \\ m \frac{\mathrm{d}^2 x_B}{\mathrm{d}t^2} = -kx_B - k(x_B - x_A) = kx_A - 2kx_B \end{cases}$$

- Consider a normal mode where both masses have the same frequency, i.e.  $\begin{cases} x_A(t) = A\cos(\omega t) \\ x_B(t) = B\cos(\omega t) \end{cases}$
- Substitute into 1:  $\frac{A}{B} = \frac{k}{2k m\omega}$ – Substitute into 2:  $\frac{A}{B} = \frac{2k - m\omega^2}{k}$ - Therefore  $2k - m\omega^2 = \pm k \implies \omega_1 = \sqrt{\frac{k}{m}}, \omega_2 = \sqrt{\frac{3k}{m}}$ \*  $\omega^2 = \frac{k}{m} \implies A = B$  – regular oscillation, middle spring inactive • General mode would be a superposition:  $\begin{cases} x_A(t) = C_1 \cos(\omega_1 t) + C_2 \cos(\omega_2 t) \\ x_A(t) = C_1 \cos(\omega_1 t) + C_2 \cos(\omega_2 t) \end{cases} \text{ where } \omega_1 = \sqrt{\frac{k}{m}}, \omega_2 = \frac{1}{2} \sum_{k=1}^{m} \frac{1}{2}$

$$\sqrt{\frac{3k}{m}}$$

- Example: Consider if  $x_A(0) = A, x_B(0) = 0$ :

  - $-q_{1} = x_{A} + x_{B}, q_{2} = x_{A} x_{B} \implies x_{A} = \frac{1}{2}(q_{1} + q_{2}), x_{B} = \frac{1}{2}(q_{1} q_{2})$   $\text{Consider } q_{1} = C_{1}\cos(\omega_{1}t), q_{2} = C_{2}\cos(\omega_{2}t); \text{ plug in } t = 0 \text{ and solve to get } C_{1} = C_{2} = A$   $x_{A} = \frac{1}{2}(C_{1}\cos(\omega_{1}t) + C_{2}\cos(\omega_{2}t)) = \frac{1}{2}A(\cos(\omega_{1}t) + \cos(\omega_{2}t)) = \frac{1}{2}A\cos\left(\frac{\omega_{1} + \omega_{2}}{2}t\right)\cos\left(\frac{\omega_{1} \omega_{2}}{2}t\right)$

\* We can think about this as a system oscillating with frequency  $\frac{\omega_1 + \omega_2}{2}$  (i.e. the average of the two), and amplitude  $\frac{\omega_1 - \omega_2}{2}t$ \* This is known as a beating phenomenon

## Extending to More Masses

• Let  $x_A = A\cos(\omega t), x_B = B\cos(\omega t)$ , substitute into the equations of motion - Recall  $\frac{\mathrm{d}^2 x}{\mathrm{d} t^2} = -\omega^2 x$ 

• 
$$\begin{cases} A(2k - m\omega^2)\cos(\omega t) = kB\cos(\omega t) \\ B(2k - m\omega^2)\cos(\omega t) = kA\cos(\omega t) \end{cases} \implies \begin{cases} \frac{2n}{m}A - \frac{n}{m}B = A\omega^2 \\ \frac{2k}{m}B - \frac{k}{m}A = B\omega^2 \end{cases}$$
  
• Write in matrix form: 
$$\begin{bmatrix} \frac{2k}{m} & -\frac{k}{m} \\ -\frac{k}{m} & \frac{2k}{m} \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} = \omega^2 \begin{bmatrix} A \\ B \end{bmatrix}$$

- This is an eigenvalue problem  $\begin{bmatrix} -\frac{m}{m} & -\frac{m}{m} \end{bmatrix}$ 
  - Solving us gets the same  $\omega$  as before
- Example: Two equal masses m suspended from identical springs of constant k each
  - Note we don't need to worry about gravity since it's only a constant offset to the equilibrium - Consider a displacement down
    - $-m\frac{\mathrm{d}^2 x_A}{\mathrm{d}t^2} = -kx_A + k(x_B x_A) \implies \frac{\mathrm{d}x_A}{\mathrm{d}t} = -\frac{2k}{m}x_A + \frac{k}{m}x_B$

$$- m \frac{\mathrm{d}^2 x_B}{\mathrm{d}t^2} = -k(x_B - x_A) = \frac{k}{m} x_A - \frac{k}{m} x_B$$
  
- Let  $x_A = A \cos(\omega t), x_B = B \cos(\omega t)$   
- As matrix equation: 
$$\begin{bmatrix} -\frac{2k}{m} & \frac{k}{m} \\ \frac{k}{m} & -\frac{k}{m} \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} = \omega^2 \begin{bmatrix} A \\ B \end{bmatrix}$$