

Lecture 32, Nov 29, 2022

Forced Convection

- Consider a plate with a boundary layer forming, with laminar, transitional, and turbulent zones, with height $\delta_v(x)$
 - The boundary layer becomes turbulent as shear stress goes down due to the velocity gradient decreasing as the boundary layer gets thicker
- Recall $\tau = C_F \frac{\rho v_\infty^2}{2}$
 - C_F , the friction coefficient, has a local value depending on where on the shape you are
 - In practice friction goes down in the laminar regime, goes up in the transition, and then goes down again the turbulent regime
 - * This is again due to the growth of the boundary layer
 - For laminar flow $C_{F,x} = \frac{0.664}{\text{Re}_x^{\frac{1}{2}}}$
 - * We can average this over the length of the plate: $C_F = \frac{1}{L} \int_0^L \frac{0.664}{\text{Re}_x^{\frac{1}{2}}} dx = \frac{1.328}{\text{Re}_L^{\frac{1}{2}}}$
 - For turbulent flow ($\text{Re}_x > 5 \times 10^5$): $C_{F,x} = \frac{0.0592}{\text{Re}_x^{\frac{1}{5}}}$
 - * This gives an average $C_F = \frac{0.074}{\text{Re}_L^{\frac{1}{5}}}$
- Now consider the thermal boundary layer
 - $h = -\frac{k}{T_s - T_\infty} \left. \frac{dT}{dy} \right|_{y=0}$
 - $\frac{dT}{dy} \sim \frac{T_s - T_\infty}{\delta_T}$
 - As the thermal boundary layer thickness goes up, the convective heat transfer goes down
 - This gives a local Nusselt number $\text{Nu}_x = \frac{h_x x}{k}$
 - * For laminar flow: $\text{Nu}_x = 0.332 \text{Re}_x^{\frac{1}{2}} \text{Pr}^{\frac{1}{3}}$ for $\text{Pr} \geq 0.6$
 - This gives an average of $\text{Nu} = 0.664 \text{Re}_L^{\frac{1}{2}} \text{Pr}^{\frac{1}{3}}$
 - * For turbulent flow: $\text{Nu}_x = 0.0296 \text{Re}_x^{\frac{4}{5}} \text{Pr}^{\frac{1}{3}}$ for $0.6 \leq \text{Pr} \leq 60, 5 \times 10^5 \leq \text{Re}_x \leq 1 \times 10^7$
 - This gives an average $\text{Nu} = 0.037 \text{Re}_L^{\frac{4}{5}} \text{Pr}^{\frac{1}{3}}$
- The film temperature is defined as the average of the surface and free temperatures $T_F = \frac{T_s + T_\infty}{2}$
 - This is to find the average fluid properties

Flows Over Cylinders and Spheres

- Behind a cylinder/sphere there is a turbulent wake, making the boundary layer behave erratically (flow separation, technically not turbulence)
 - Turbulence can still occur, if $\text{Re} = \frac{vD}{\nu} > 2 \times 10^5$
- For a cylinder $\text{Nu} = C \text{Re}^m \text{Pr}^n$ just like for a plate
- The Churchill and Bernstein correlation is more accurate and applies over a broader range
 - $\text{Nu} = 0.3 + \frac{0.62 \text{Re}^{\frac{1}{2}} \text{Pr}^{\frac{1}{3}}}{\left(1 + \left(\frac{0.4}{\text{Pr}}\right)^{\frac{2}{3}}\right)^{\frac{1}{4}}} \left(1 + \left(\frac{\text{Re}}{28200}\right)^{\frac{5}{8}}\right)^{\frac{4}{5}}$, applicable over $\text{RePr} > 0.2$
- Flow over a sphere is similar: $\text{Nu} = 2 + \left(0.4 \text{Re}^{\frac{1}{2}} + 0.06 \text{Re}^{\frac{2}{3}}\right) \text{Pr}^{0.4} \left(\frac{\mu_\infty}{\mu_s}\right)^{\frac{1}{4}}$, valid for $3.5 \leq \text{Re} \leq 80000, 0.7 \leq \text{Pr} \leq 380$
 - Note all properties are evaluated at T_∞ , not T_p like the others (except μ_s)

Summary

For laminar flow over a flat plate (average):

$$\text{Nu} = 0.664\text{Re}_L^{\frac{1}{2}}\text{Pr}^{\frac{1}{3}}, \text{Pr} \geq 0.6$$

For turbulent flow over a flat plate (average):

$$\text{Nu} = 0.037\text{Re}_L^{\frac{4}{5}}\text{Pr}^{\frac{1}{3}}, 0.6 \leq \text{Pr} \leq 60, 5 \times 10^5 \leq \text{Re} \leq 1 \times 10^7$$

where Re is the Reynolds number evaluated for the entire plate; material properties and Pr can be determined through the film temperature

$$T_F = \frac{T_s + T_\infty}{2}$$