Lecture 27, Nov 17, 2022

Fin Sizing

- When can we assume a fin to be infinitely long?
- Compare adiabatic tip vs. infinitely long fin equation, taking the ratio we get tanh(aL)
- For aL = 1, tanh(aL) = 0.762; for aL = 5, tanh(aL) = 0.9999)
 - As a rule of thumb if $aL \geq 5$ we can assume the fin is infinitely long
 - A value of 1 gives 76.2% of the total heat transfer from an infinitely long fin, with a lot less material
- $L = \frac{1}{a}$ is typically a reasonable length for a fin
- What about the area of fins?
 - Consider area with fin and area without fin
 - $\dot{Q}_{total} = \dot{Q}_{nofin} + \dot{Q}_{fin}$
 - From the fin efficiency definition $\eta_{fin} = \frac{\dot{Q}_{fin}}{hA_{fin}(T_b T_\infty)}$

- $\dot{Q}_{fin} = h\eta_{fin}A_{fin}(T_b - T_\infty)$ - Putting it all together $\dot{Q}_{total} = h(A_{nofin} + \eta_{fin}A_{fin})(T_b - T_\infty)$

- $(A_{nofin} + \eta_{fin}A_{fin})$ is the effective area for heat transfer In terms of thermal resistances $R = \frac{T_b T_{\infty}}{\dot{Q}_{total}} = \frac{1}{h(A_{nofin} + \eta_{fin}A_{fin})}$

Transient Heat Conduction (Lumped)

• Simple example: taking a material and immersing it in a fluid with a high temperature difference, resulting in rapid heat transfer

- What is
$$T(t)$$
?

- Recall $\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} = \frac{1}{\alpha} \frac{\partial T}{\partial t}$
- Lumped capacitance: simple assumption that there is no temperature gradient in the body, i.e. $\frac{\partial^2 T}{\partial m^2} =$

 $\frac{\partial^2 T}{\partial u^2} = \frac{\partial^2 T}{\partial z^2} = 0$, the temperature is at uniform temperature everywhere

- In reality modes of heat transfer include conduction within the material and convection to the surrounding fluid
- This makes sense only if $R_{cond} \ll R_{conv}$

• Energy balance: Let $\Delta \dot{E}(t) = -\dot{Q}_{conv}(t)$ be the energy change in the body

$$-\Delta \dot{E}(t) = mc_p \frac{dT}{dt} = \rho V c_p \frac{dT}{dt}$$
$$-hA(T - T_{\infty}) = -\rho V c_p \frac{dT}{dt} \implies \frac{dT}{T - T_{\infty}} = \frac{d(T - T_{\infty})}{T - T_{\infty}} = -\frac{hA}{\rho V c_p} dt$$

- Integrate to get $\ln(T - T_{\infty}) = -\frac{nA}{\rho V c_n} t + c_1$, apply boundary condition that $T(0) = T_i$

$$-\ln\left(\frac{T-T_{\infty}}{T_i-T_{\infty}}\right) = -\frac{hA}{\rho V c_p}$$

- This compares the "changing driving force" against the "max driving force" • Define the time constant $\tau = \frac{\rho V c_p}{hA}$ so $\frac{T - T_{\infty}}{T_i - T_{\infty}} = e^{-\frac{t}{\tau}}$

Equation

For lumped transient heat conduction,

$$\frac{T - T_{\infty}}{T_i - T_{\infty}} = e^{-\frac{t}{\tau}}$$

where $\tau = \frac{\rho V c_p}{hA}$

Validity of the Lumped Capacitance Assumption

- This only makes sense if $R_{conv} \gg R_{cond}$
- Take $\frac{R_{cond}}{R_{conv}} = \frac{\frac{L}{kA}}{\frac{1}{hA}} = \frac{hL}{k}$
- Bi = $\frac{hL_c}{k}$ is the *Biot number*, a dimensionless quantity

- L_c^{κ} is a characteristic length in the direction of conduction, from the midpoint to the wall - $L_c = \frac{V}{C}$

$$L_c = \overline{A}$$

- For a sphere $L_c = \frac{r}{3}$, for a cylinder (with length \gg radius) $L_c = \frac{r}{2}$

- Consider the steady state analogue; for $Bi \gg 1$, the temperature drops the sharpest over conduction, so lumped capacitance is not valid; for $Bi \ll 1$, the temperature drops the sharpest over convection, so lumped capacitance is valid
- Typical cutoff is Bi < 0.1•
- Example: putting steel rod at 300°C into furnace at 1200°C with $h = 100 \text{W/m}^2 \text{ K}, D = 0.1m, k =$

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- 0.1 so the assumption is valid Using lumped capacitance $\frac{T T_{\infty}}{T_i T_{\infty}} = e^{-\frac{hA}{\rho V_c}t} = e^{-\frac{h}{\rho c} \cdot \frac{2}{r}t}$ $\ln \frac{800 1200}{300 1200} = \frac{-2 \cdot 100 \text{W}m^2.K}{7832 \text{kg/m}^3 \cdot 541 \text{J}kg.K}t$

$$-t = 859s$$