## Lecture 26, Nov 15, 2022

## More On Finned Surfaces

- Why does longer hair not lead to more heat transfer?
  - -k for hair is very low, so longer hair leads to a negligible amount of additional heat transfer
  - The additional hair creates a boundary layer that effectively lowers h leading to less convection
- Consider cylindrical fins, increasing the diameter from d to 2d results in  $\frac{Q'}{\dot{Q}} = 2^{\frac{3}{2}}$ , an increase in the

total heat transfer

- For heat flux however we get  $\frac{\dot{q}'}{\dot{q}} = 2^{-\frac{1}{2}}$ , which is lower
- When designing a heat sink it might be better to have a larger number of smaller fins
  - \* Note: Having fins that are too small might break up the flow and change h, leading to worse performance

## **Finite Length Fins**

• Method 1: Consider an adiabatic tip (insulated tip)

= 0

- No heat transfer at the tip means the temperature at the tip must be constant

$$-\frac{\mathrm{d}T}{\mathrm{d}t} = \frac{\mathrm{d}\theta}{\mathrm{d}t}$$

$$dx \mid_{x=L} dx \mid_{x=L}$$

- $-\theta(x) = c_1 e^{ax} + c_2 e^{-ax} \implies \frac{\mathrm{d}\theta}{\mathrm{d}x} = c_1 a e^{ax} c_2 a e^{-ax}$  Plug in the initial condition,  $0 = c_1 a e^{aL} c_1 a e^{-aL}$

- At  $x = 0, \theta = \theta_b$ , so  $\theta_b = c_1 + c_2$  Solving yields  $\frac{\theta(x)}{\theta_b} = \frac{\cosh(a(L-x))}{\cosh(aL)}$
- Solving for  $\dot{Q}$  using Fourier's law yields  $\sqrt{hPkA_c}(T_s T_\infty) \tanh(aL)$ 
  - \* Note  $tanh(L) \to 1$  as  $L \to \infty$ , so this approaches the infinitely long fin equation as the fin gets longer
- Method 2: Use the "corrected length"
  - Have the convection coming out of the tip be idealized as coming out from the fin side
  - Imagine extending the fin by  $\Delta L$  such that the additional side area  $\Delta LP$  equals the size of the tip cross-section, now we can assume the tip is adiabatic
  - Corrected length is  $L_c = L + \frac{A_c}{P}$
  - Use this  $L_c$  with the adiabatic tip solution for  $T, \theta$
  - Note this only works well if  $A_c \ll L$  so the extension is minimal

Summary

For a fin with finite length L and an insulated tip, then

$$\frac{\theta(x)}{\theta_b} = \frac{T(x) - T_{\infty}}{T_b - T_{\infty}} = \frac{\cos(a(L-x))}{\cosh(aL)}$$

where

$$a = \sqrt{\frac{hP}{kA_c}}$$

and the total heat transfer through the fin is

$$\dot{Q} = \sqrt{hPkA_c(T_s - T_\infty) \tanh(aL)}$$

For fins with cross sectional area small relative to the length and non-adiabatic tip, use the corrected length  $L_c = L + \frac{A_c}{P}$ 

## **Fin Efficiency**

- The most "efficient" fin would have effectively infinite conductivity, so  $T(x) = T_b$  for all x so that the convection along the fin is maximized
- In this case the heat transfer is just convection at uniform temperature,  $Q_{max} = hA_{fin}(T_b T_\infty) = hPL(T_b T_\infty)$  ignoring the fin tip
- Define the fin efficiency as  $\eta_{fin} = \frac{\dot{Q}}{\dot{Q}_{max}}$ - This is equal to  $\frac{\sqrt{hPkA_c}(T_b - T_\infty)}{hPL(T_b - T_\infty)} = \frac{1}{aL}$  for an infinitely long fin - For an adiabatic tip  $\frac{\sqrt{hPkA_c}(T_b - T_\infty) \tanh(aL)}{hPL(T_b - T_\infty)} = \frac{\tanh(aL)}{aL}$ - As L increases,  $\eta_{fin}$  approaches 0
- The fin effectiveness is defined as  $\varepsilon_{fin} = \frac{\dot{Q}_{fin}}{\dot{Q}_0}$  where  $\dot{Q}_0$  is the heat transfer without the fin, which would be  $hA_c(T_b T_\infty)$ 
  - For an infinitely long fin  $\varepsilon_{fin} = \sqrt{\frac{kP}{hA_c}}$
- To increase the fin effectiveness, maximize k and  $\frac{P}{A_c}$
- When h goes up, the fin effectiveness goes down; fins are the most effective with low h
  e.g. if we have a boundary between air and water, it's better to have the fin on the air side since h in air is much lower
- Rule of thumb: Fins are worth it if  $\varepsilon \geq 2$