

## Lecture 21, Oct 27, 2022

### One Dimensional Heat Conduction Equation

- Consider heat conduction  $x \rightarrow x + \Delta x$ , surface area at  $x$  is  $A_x$ ; what is the temperature as a function of  $x$ ?
- $\frac{dh}{dt} = mc_p \frac{\partial T}{\partial t} = \rho V c_p \frac{\partial T}{\partial t} = \rho c_p A \Delta x \frac{\partial T}{\partial t}$
- Energy balance:  $\rho c_p A \Delta x \frac{\partial T}{\partial t} = \dot{Q}_x - \dot{Q}_{x+\Delta x}$ 
  - In terms of heat flux,  $\dot{q}_x A_x - \dot{q}_{x+\Delta x} A_{x+\Delta x}$
  - $\rho c_p \frac{\partial T}{\partial t} = -\frac{1}{A} \left( \frac{\dot{q}_x A_x - \dot{q}_{x+\Delta x} A_{x+\Delta x}}{\Delta x} \right)$
  - Take the limit  $\Delta x \rightarrow 0$ :  $\rho c_p \frac{\partial T}{\partial t} = -\frac{1}{A} \frac{\partial(\dot{q}A)}{\partial x}$

### Cartesian Coordinates

- Consider Cartesian coordinates, constant area, then we can simplify this as  $\rho c_p \frac{\partial T}{\partial t} = -\frac{\partial \dot{q}}{\partial x}$
- Putting this into Fourier's law,  $\dot{q} = -k \frac{dT}{dx}$ , we get  $\rho c_p \frac{\partial T}{\partial t} = -\frac{\partial}{\partial x} \left( -k \frac{\partial T}{\partial x} \right)$ 
  - With a constant  $k$ , we get  $\rho c_p \frac{\partial T}{\partial t} = k \frac{\partial^2 T}{\partial x^2}$
  - Alternatively  $\frac{\partial T}{\partial t} = \frac{k}{\rho c_p} \frac{\partial^2 T}{\partial x^2} = \alpha \frac{\partial^2 T}{\partial x^2}$
- $\alpha = \frac{k}{\rho c_p}$  is the thermal diffusivity, with units of  $m^2/s$ 
  - Higher  $k$  conducts heat well so the gradient is sharper
  - $\rho c_p$  stores energy well, so a lot of heat can enter the system without changing the temperature much
  - $\alpha = 1.11 \times 10^{-4} m^2/s$  for copper,  $\alpha = 3.4 \times 10^{-7} m^2/s$
- For steady state,  $\frac{\partial T}{\partial t} = 0$  so  $\alpha \frac{\partial^2 T}{\partial x^2} = 0$ 
  - Integrate this and we get that  $\frac{\partial T}{\partial x}$  is a constant

### Cylindrical Coordinates

- In the radial direction Fourier's law is  $\dot{q} = -k \frac{\partial T}{\partial r}$
- $A = 2\pi r L$
- $\rho c_p \frac{\partial T}{\partial t} = -\frac{1}{2\pi r L} \left( \frac{\partial}{\partial r} \left( 2\pi r L \left( -k \frac{\partial T}{\partial r} \right) \right) \right) = \frac{k}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right)$
- $\frac{\partial T}{\partial t} = \frac{\alpha}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right)$
- For steady state, this simplifies to  $\frac{\alpha}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right) = 0 \implies r \frac{\partial T}{\partial r} = 0$

### Spherical Coordinates

- $A = 4\pi r^2$ ,  $\dot{q} = -k \frac{\partial T}{\partial r}$
- Doing the same derivation gets us  $\rho c_p \frac{\partial T}{\partial t} = \frac{k}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right)$
- $\frac{\partial T}{\partial t} = \frac{\alpha}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right)$

## Summary

1D heat flow equations:

- Cartesian coordinates:

$$\frac{\partial T}{\partial t} = \alpha \frac{\partial^2 T}{\partial x^2}$$

- Cylindrical coordinates (radial):

$$\frac{\partial T}{\partial t} = \frac{\alpha}{r} \frac{\partial}{\partial r} \left( r \frac{\partial T}{\partial r} \right)$$

- Spherical coordinates (radial):

$$\frac{\partial T}{\partial t} = \frac{\alpha}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial T}{\partial r} \right)$$

where  $\alpha$  is the thermal diffusivity,  $\alpha = \frac{k}{\rho c_p}$

In general,

$$\frac{\partial T}{\partial t} = \frac{\alpha}{r^n} \frac{\partial}{\partial r} \left( r^n \frac{\partial T}{\partial r} \right)$$

where  $n = 0$  for Cartesian,  $n = 1$  for cylindrical and  $n = 2$  for spherical