

Lecture 31, Dec 6, 2021

Simple Harmonic Oscillator Energy

- Recall $a(t) = -A\omega^2 \cos(\omega t + \phi) \implies F_x = -mA\omega^2 \cos(\omega t + \phi) = -m\omega^2 x(t)$
- We can integrate this force to find the work done: $W = \int_{x_0}^x -m\omega^2 x \, dx = \frac{1}{2}m\omega^2(x_0^2 - x^2)$
- $W = \Delta K \implies \Delta K = \frac{1}{2}m\omega^2(x_0^2 - x^2)$
- If we make the oscillator a closed system then $\Delta U = -\Delta K = \frac{1}{2}m\omega^2(x^2 - x_0^2)$
 - Since potential at x_0 is arbitrary we can set this to 0 so $\Delta U = \frac{1}{2}m\omega^2 x^2$
- The total energy is then $E = U + K = \frac{1}{2}m\omega^2 x^2 + \frac{1}{2}mv^2 = \frac{1}{2}m\omega^2 A^2 \cos^2(\omega t + \phi) + \frac{1}{2}m\omega^2 A^2 \sin^2(\omega t + \phi) = \frac{1}{2}m\omega^2 A^2$
 - The energy of a simple harmonic oscillator is constant and it only trades potential and kinetic energy back and forth
 - The energy is proportional to the square of the amplitude

Torsional Oscillators

- A disk suspended by a twisting fibre creates a torsional oscillator in simple harmonic motion
 - $\tau = I\alpha$
 - For small angular displacements $\tau = -\kappa(\omega - \omega_0)$ where κ is the angular equivalent of k
 - The differential equation is then $\frac{d^2\theta}{dt^2} = -\frac{\kappa}{I}\theta$ which is the same as the one for translational simple harmonic motion, just with the translational terms substituted by rotational ones
- The equations of motion are identical to translational simple harmonic motion $\theta = \theta_{max} \cos(\omega t + \phi)$ and $\omega = \sqrt{\frac{\kappa}{I}}$

Examples of Oscillating Systems

- All simple harmonic oscillators obey $\frac{d^2x}{dt^2} + \omega^2 x = 0$
- Examples:
 1. Mass on a spring: $m \frac{d^2x}{dt^2} + kx = 0 \implies \omega^2 = \frac{k}{m}$
 2. Torsional oscillator: $I \frac{d^2\theta}{dt^2} + \kappa\theta = 0 \implies \omega^2 = \frac{\kappa}{I}$
 3. Pendulum: $ml^2 \frac{d^2\theta}{dt^2} + mgl \sin \theta \approx ml^2 \frac{d^2\theta}{dt^2} + mgl\theta = 0 \implies \omega^2 = \frac{g}{l}$
 4. Floating object bobbing in water: $m \frac{d^2y}{dt^2} + g(\rho Ay) = 0 \implies \omega^2 = \frac{g\rho A}{m}$
 - A buoyant object will float in the water at some neutral point, and if pushed past this neutral point then $F_y \propto y$
 5. Capacitor-inductor circuit $L \frac{d^2Q}{dt^2} + \frac{Q}{C} = 0 \implies \omega^2 = \frac{1}{LC}$
 - Voltage drop across capacitor: $\frac{Q}{C}$
 - Voltage drop across inductor: $L \frac{d^2Q}{dt^2}$

Damped Oscillations

- In reality there is always some friction present, causing the oscillator to lose energy and thus amplitude

- The loss in amplitude is called *damping*
- Often the damping is caused by viscous friction, with the damping force opposing and proportional to velocity $F = -b\vec{v}$
- Therefore $F_x = -(kx + bv_x) = ma_x \implies m\frac{d^2x}{dt^2} + b\frac{dx}{dt} + kx = 0$
- The solution is given by $x(t) = Ae^{-\frac{bt}{2m}} \cos(\omega_d t + \phi)$ where $\omega_d = \sqrt{\omega^2 - \left(\frac{b}{2m}\right)^2} = \sqrt{\frac{k}{m} - \frac{b^2}{4m^2}}$
 - A damped harmonic oscillator oscillates slower than the equivalent undamped oscillator
 - Let the damping time constant $\gamma = \frac{b}{m}$, with units of time
 - The amplitude can be expressed as $x_{max} = Ae^{-\frac{\gamma t}{2}}$
 - The energy remaining is then $E(t) = \frac{1}{2}m\omega^2 x_{max}^2 = \left(\frac{1}{2}m\omega^2 A^2\right) e^{-\gamma t}$
- As b increases the oscillator decays faster, and when $b > 2m\omega$ the system is overdamped and there are no oscillations at all

Q Factor

- Q is the quality factor of an oscillator and measures the rate of decay
 - Differs from ω in the sense that Q measures the number of oscillations
- $Q \equiv \frac{\omega}{\gamma} = \frac{2\pi}{\gamma T}$
- If $Q = 2\pi$ then the energy falls to $e^{-1} = 37\%$ of its original energy in a single cycle
- A good bell has $Q = 100$, electronic circuits have $Q = 10^6$, quantum systems have $Q = 10^9$